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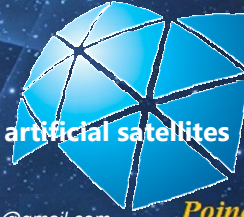
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Semi-analytical study of the stability of the rotational motion of artificial satellites subject to external torques using quaternions

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Poincaré
(1854–1912)



Introduction

This study aims to perform the stability analysis of the rotational motion to artificial satellites using quaternions to describe the satellite attitude (orientation on the space). In the system of rotational motion equations, which is composed by the four kinematic equations of the quaternions and by the three Euler equations in terms of the rotational spin components. The torques that we have been considered are the Gravity Gradient (GGT) and The Direct Solar Radiation Pressure (DSRPT). Equilibrium points are obtained through numerical simulations using the softwares MATLAB and OCTAVE, which are then analyzed by the Routh-Hurwitz Stability Criterion.

Attitude Quaternions and Rotational Motion Equations

Attitude Quaternions [1]

$$\hat{n} = (n_1 \ n_2 \ n_3)^t \quad (1) \quad q_1 = n_1 \operatorname{sen}\left(\frac{\phi}{2}\right) \quad q_2 = n_2 \operatorname{sen}\left(\frac{\phi}{2}\right) \quad (2)$$

$$q = (\bar{q} \ q_4)^t = (q_1 \ q_2 \ q_3 \ q_4)^t \quad q_3 = n_3 \operatorname{sen}\left(\frac{\phi}{2}\right) \quad q_4 = \cos\left(\frac{\phi}{2}\right)$$

Motion Equations

Kinematic Equations [2]

$$\begin{aligned} \dot{q}_1 &= 0.5[pq_4 - qq_3 + rq_2] \\ \dot{q}_2 &= 0.5[qq_4 - rq_1 + pq_3] \\ \dot{q}_3 &= 0.5[rq_4 - pq_2 + qq_1] \\ \dot{q}_4 &= -0.5[pq_1 + qq_2 + rq_3] \end{aligned} \quad (3)$$

Euler Equations [3]

$$\begin{aligned} I_x \dot{p} &= M_x + (I_y - I_z) q r \\ I_y \dot{q} &= M_y + (I_z - I_x) p r \\ I_z \dot{r} &= M_z + (I_x - I_y) p q \end{aligned} \quad (4)$$

Gravity Gradient Torque

Gravity Gradient Elementar Force over an infinitesimal area element (dS) of the satellite surface, is given by [1, 2]:

$$d\vec{F} = -\mu dm \frac{\vec{r}}{r^3} \quad (5)$$

Gravity Gradient Torque over an infinitesimal area element (dS) of the satellite surface, is given by:

$$d\vec{N}_G = \vec{r}_G \times d\vec{F} \quad \vec{N}_G = 3 \frac{\mu}{r^3} [a_x a_x (I_z - I_x) \hat{i}_x + a_x a_x (I_z - I_x) \hat{j}_x + a_x a_x (I_z - I_x) \hat{k}_x] \quad (6)$$

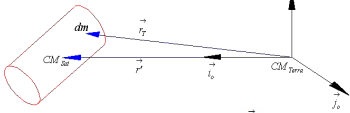


Fig. 1: Relations between the mass elements of the satellite and the Earth, being \vec{r} the satellite center of mass position vector related to the Earth center of mass, and \hat{r} the unit vector in its direction.

Solar Radiation Torque (SRT)

Solar Radiation Pressure Elementar Force over an infinitesimal area element (dS) of the satellite surface [4]:

$$d\vec{F} = -\frac{K}{R^2} \left[\frac{2\gamma}{3} (1 - \beta) \cos\theta_p + 4\beta\gamma \cos^2\theta_p \right] \hat{u} + [(1 - \beta\gamma) \cos\theta_p] \hat{u} \quad (7)$$

Direct Solar Radiation Pressure Torque over an infinitesimal area element (dS) of the satellite surface:

$$d\vec{N} = \vec{r} \times d\vec{F} \quad (8)$$

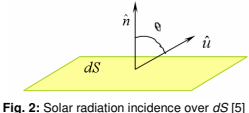


Fig. 2: Solar radiation incidence over dS [5]

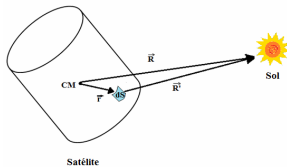


Fig. 3: CM and dS positions related to the Sun [6]

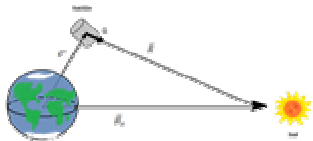


Fig. 4: Vectors arrangement connecting the satellite, the Earth and the Sun [6]

$$\vec{N} = N_x \hat{e}_x + N_y \hat{e}_y + N_z \hat{e}_z \quad (9)$$

$$N_x = -\frac{K}{R^4} (\beta_1 \gamma_1 - \beta_2 \gamma_2) \frac{h}{2} \pi \sigma^2 G_1 \quad N_y = -\frac{K}{R^4} (\beta_1 \gamma_1 - \beta_2 \gamma_2) \frac{h}{2} \pi \sigma^2 G_2 \quad N_z = 0 \quad (10)$$

Algorithm for Stability Analysis

The Routh-Hurwitz Criterion allows to investigate the absolute stability of a system of equations using the coefficients of the characteristic equation associated with the linearized system, without the need to determine the roots of the characteristic equation. [7, 8].

$$a_n \lambda^n + a_{n-1} \lambda^{n-1} + \dots + a_2 \lambda^2 + a_1 \lambda + a_0 = 0 \quad (11)$$

Routh Table, general case:

k \ j	1	2	3	...	j
k = i - 1					
0	$\lambda^n : a_n$	a_{n-2}	a_{n-4}	...	a_{n-2i+2}
1	$\lambda^{n-1} : a_{n-1}$	a_{n-3}	a_{n-5}	...	$a_{n-1-2i+2}$
2	$\lambda^{n-2} : b_{n-2}$	b_{n-4}	b_{n-6}	...	$b_{n-2-2i+2}$
3	$\lambda^{n-3} : c_{n-3}$	c_{n-5}	c_{n-7}	...	$c_{n-3-2i+2}$
k	$\lambda^{n-k} : Z_{n-k}$	Z_{n-k-2}	Z_{n-k-4}	...	$Z_{n-k-2-2i+2}$
n - 4	$\lambda^4 : \alpha_4$	α_2	α_0		
n - 3	$\lambda^3 : \beta_3$	β_1			
n - 2	$\lambda^2 : \gamma_2$	γ_0			
n - 1	$\lambda^1 : \delta_1$				
n	$\lambda^0 : \epsilon_0$				

Routh Table, case n = 7:

$$\begin{aligned} &\lambda^7 \ a_7 \ a_5 \ a_3 \ a_1 \\ &\lambda^6 \ a_6 \ a_4 \ a_2 \ a_0 \\ &\lambda^5 \ b_5 \ b_3 \ b_1 \\ &\lambda^4 \ c_4 \ c_2 \ c_0 \\ &\lambda^3 \ d_3 \ d_1 \ e_0 \\ &\lambda^2 \ e_2 \ e_0 \\ &\lambda^1 \ f_1 \\ &\lambda^0 \ g_0 \end{aligned}$$

Results obtained by numerical simulations

Table 1: Order of Magnitude of the eigenvalues with real part less than 10^{-5}

Satellite	10^{-5}	10^{-6}	10^{-7}	10^{-8}	10^{-9}	10^{-10}	10^{-14}	TOTAL
MP	9	2	2	0	1	0	0	14
PP-1	36	0	10	2	4	2	2	56
PP-2	39	1	3	1	3	1	1	49
TOTAL	85	3	15	3	7	3	2	119

Table 2: Satellite PP-2 equilibrium point

Equilibrium point	Values obtained in the
p (rad/s)	4.239×10^{-13}
q (rad/s)	1.578×10^{-13}
r (rad/s)	-9.198×10^{-13}
q_1	9.001×10^{-2}
q_2	9.001×10^{-2}
q_3	9.001×10^{-1}
q_4	6.079×10^{-12}

Table 3: 1st Column of the Routh Table for a PP-2 equilibrium point

Coefficients of the first column	Values obtained in the simulations
a_0	1.000
a_1	7.109×10^{-20}
b_1	-1.208
c_1	8.577×10^{-20}
d_1	-1.995×10^{-14}
e_1	1.340×10^{-28}
f_1	1.481×10^{-22}
g_1	1.455×10^{-38}

Table 4: Eigenvalues associated to the PP-2 instable

Eigenvalues for the instable point	Real Part	Imaginary Part
1	8.697×10^{-14}	6.528×10^{-4}
2	8.697×10^{-14}	-6.528×10^{-4}
3	6.500×10^{-14}	3.108×10^{-4}
4	6.500×10^{-14}	-3.108×10^{-4}
5	-1.482×10^{-13}	6.101×10^{-5}
6	-1.482×10^{-13}	-6.101×10^{-5}
7	-9.489×10^{-15}	0

Conclusions

- The main significance of this work is due to the use of quaternions in determining the SRT components and in a preliminary analysis of stability. Previous works used Andoyer variables and Euler angles.
- The approach taken can be improved when in the stability analysis are included other environmental torques as a mean to obtain a better representation of the real behavior of the satellite, and also to introduce the Shadow Function in the DSRPT model.
- Other methods of stability analysis can also be used as the Method, the Liapunov Direct Method and the Kovalev Savchenko Stability Theorem [9,10].
- The latter method requires the determination of the equations of motion in its canonical form, and it is therefore also a challenge to obtain the canonical form for these equations of the quaternions.
- The approach taken here is applied to satellites with cylindrical form in the illuminated phase of the trajectory. It can be useful for analyzing the stability of future missions of the

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